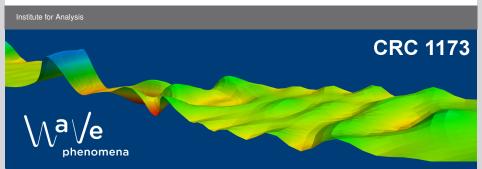


Localized time-periodic solutions of nonlinear wave equations

Wolfgang Reichel (financially supported by DFG through CRC 1173)

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• May 22–26, 2017



The problem



Find spatially localized, time-periodic $E: \mathbb{R}^3 \times \mathbb{R} \to \mathbb{R}^3$ such that

(quasi)
$$\nabla \times \nabla \times E + \partial_t^2 (V(x)E + \Gamma(x)|E|^{p-1}E) = 0$$

(semi)
$$\nabla \times \nabla \times E + V(x)\partial_t^2 E + \Gamma(x)|E|^{p-1}E = 0$$

with p > 1 & suitable conditions on $V, \Gamma : \mathbb{R}^3 \to \mathbb{R}$

Outline:

- (A) physical background
- (B) results for semilinear wave equation
- (C) results for quasilinear wave equation



$$\nabla \times E + \partial_t B = 0$$
,

$$\nabla \cdot D = 0$$
,

$$\nabla \times H - \partial_t D = 0,$$

$$\nabla \cdot B = 0$$
.

Material laws:

$$B = \mu_0 H$$
, $D = \epsilon_0 E + P(x, E) = \epsilon_0 (1 + \chi_1(x) + \chi_3(x) |E|^2 + ...) E$



$$\nabla \times E + \frac{\partial_t B}{\partial_t B} = 0$$
,

$$\nabla \cdot D = 0$$
,

$$\nabla \times \mathbf{H} - \partial_t \mathbf{D} = \mathbf{0},$$

$$\nabla \cdot B = 0$$
.

Material laws:

$$B = \mu_0 H$$
, $D = \epsilon_0 E + P(x, E) = \epsilon_0 (1 + \chi_1(x) + \chi_3(x)|E|^2 + ...)E$

$$\hookrightarrow$$

$$\nabla \times \nabla \times E + \partial_t^2(\mu_0 D) = 0$$



$$\nabla \times E + \partial_t B = 0$$
,

$$\nabla \cdot D = 0$$
,

$$\nabla \times H - \partial_t D = 0,$$

$$\nabla \cdot B = 0$$
.

Material laws:

$$B = \mu_0 H$$
, $D = \epsilon_0 E + P(x, E) = \epsilon_0 (1 + \chi_1(x) + \chi_3(x)|E|^2 + ...)E$

$$\hookrightarrow$$

$$\nabla \times \nabla \times E + \partial_t^2 \left(\mu_0 \mathbf{D} \right) = 0$$



$$\nabla \times E + \partial_t B = 0$$
,

$$\nabla \cdot D = 0$$
,

$$\nabla \times H - \partial_t D = 0,$$

$$\nabla \cdot B = 0$$
.

Material laws:

$$B = \mu_0 H$$
, $D = \epsilon_0 E + P(x, E) = \epsilon_0 (1 + \chi_1(x) + \chi_3(x)|E|^2 + ...)E$

$$\hookrightarrow$$

$$\nabla \times \nabla \times E + \partial_t^2 \left(\mu_0 \mathbf{D} \right) = 0$$

Quasilinear wave-equation:

$$\hookrightarrow \nabla \times \nabla \times E + \partial_t^2 \left(\underbrace{\mu_0 \epsilon_0 (1 + \chi_1(x))}_{=V(x) \ge 0} E + \underbrace{\mu_0 \epsilon_0 \chi_3(x) |E|^2 E + \ldots}_{=f(x,|E|^2) E} \right) = 0$$



$$\nabla \times E + \partial_t B = 0$$
,

$$\nabla \cdot D = 0$$
,

$$\nabla \times H - \partial_t D = 0,$$

$$\nabla \cdot B = 0$$
.

Material laws:

$$B = \mu_0 H$$
, $D = \epsilon_0 E + P(x, E) = \epsilon_0 (1 + \chi_1(x) + \chi_3(x)|E|^2 + ...)E$

$$\hookrightarrow$$

$$\nabla \times \nabla \times E + \partial_t^2 \left(\mu_0 \mathbf{D} \right) = 0$$

Quasilinear wave-equation:

$$\hookrightarrow \nabla \times \nabla \times E + \partial_t^2 \underbrace{\left(\underline{\mu_0 \epsilon_0 (1 + \chi_1(x))}_{=V(x) \ge 0} E + \underline{\mu_0 \epsilon_0 \chi_3(x) |E|^2 E + \ldots}\right)}_{=f(x,|E|^2)E} = 0$$

Semilinear (approximative/toy-model) variant:

$$\hookrightarrow$$

$$\nabla \times \nabla \times E + V(x)\partial_t^2 E + f(x, |E|^2)E = 0$$

(B): Semilinear wave-equations



Find solutions $U: \mathbb{R}^3 \times \mathbb{R} \to \mathbb{R}^3$ such that

(*)
$$\begin{cases} \nabla \times \nabla \times U + V(x)U_{tt} + f(x,|U|^2)U & = & 0 \\ U(x,t) & \to & 0 \text{ as } |x| \to \infty \\ U(x,t+T) & = & U(x,t) \end{cases}$$

under suitable conditions on V, f.

U (real-valued, time-periodic & spatially localized) is called "breather"

Motivation:

The famous Sine-Gordon breather

$$u_{tt} - u_{xx} + \sin u = 0$$

 The example by Blank, Chirilus-Bruckner, Lescaret, Schneider ('11) for

$$V(x)u_{tt} - u_{xx} + q(x)u = \Gamma(x)u^3$$

Source: Wikipedia

$$u(x,t) = 4 \arctan\left(\frac{m \sin(\omega t)}{\omega \cosh(mx)}\right)$$

 $m^2 + \omega^2 = 1$

(B) A vector-valued breather example



$$(*) \quad V(x)\partial_t^2 U + \nabla \times \nabla \times U + q(x)U \pm \Gamma(x)|U|^{p-1}U = 0 \quad \text{in} \quad \mathbb{R}^3 \times \mathbb{R}$$

ansatz: $U(x, t) = \psi(r, t) \frac{x}{r}, \quad r = |x|.$

(B) A vector-valued breather example



(*)
$$V(x)\partial_t^2 U + \nabla \times \nabla \times U + q(x)U \pm \Gamma(x)|U|^{p-1}U = 0$$
 in $\mathbb{R}^3 \times \mathbb{R}$

ansatz: $U(x,t) = \psi(r,t)\frac{x}{r}, \quad r = |x|.$

Theorem 1 (Plum, R. JEPE 2017)

Let
$$T = 2\pi \sqrt{\frac{V(0)}{q(0)}}$$
.

- \lor V, q, Γ > 0 radially symmetric \mathbb{C}^2 -functions,
- \blacksquare sup $\frac{q}{r} < \infty$,

$$\left|\frac{q(r)}{V(r)} - \frac{q(0)}{V(0)}\right|^{\frac{1}{p-1}} = \left\{\begin{array}{l} O(e^{-\alpha r}) \text{ as } r \to \infty, \\ o(1) \text{ in } C^2\text{-sense as } r \to 0. \end{array}\right.$$

Then \exists T-periodic, real-valued, exponentially decaying solution.



(B) The proof – solving an ODE
$$U(r,t)=\psi(r,t)\frac{x}{r}, \qquad V(r)\ddot{\psi}+q(r)\psi\pm\Gamma(r)|\psi|^{p-1}\psi=0$$

ODE in time with r = parameter

(B) The proof – solving an ODE

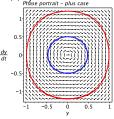


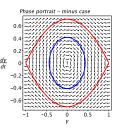
$$U(r,t) = \psi(r,t)\frac{x}{r}, \qquad V(r)\ddot{\psi} + q(r)\psi \pm \Gamma(r)|\psi|^{p-1}\psi = 0$$

ODE in time with r =parameter

Rescale:
$$\psi(r, t) = \left(\frac{q(r)}{\Gamma(r)}\right)^{1/(p-1)} y\left(\sqrt{\frac{q(r)}{V(r)}}t\right)$$

$$\ddot{y} + y \pm |y|^{p-1}y = 0$$





(B) The proof – solving an ODE



$$U(r,t) = \psi(r,t)\frac{x}{r}, \qquad V(r)\ddot{\psi} + q(r)\psi \pm \Gamma(r)|\psi|^{p-1}\psi = 0$$

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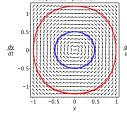
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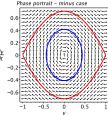
$$\ddot{y} + y \pm |y|^{p-1}y = 0$$

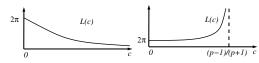
periodic orbits y(t; c), period L(c)

- c= value of first integral
- How to choose c = c(r)?

• Answer:
$$\underbrace{\sqrt{\frac{q(r)}{V(r)}T}}_{\leq 2\pi} = L(c)$$







(B) The proof – solving an ODE



$$U(r,t) = \psi(r,t)\frac{x}{r}, \qquad V(r)\ddot{\psi} + q(r)\psi \pm \Gamma(r)|\psi|^{p-1}\psi = 0$$

ODE in time with r = parameter

Rescale:
$$\psi(r,t) = \left(\frac{q(r)}{\Gamma(r)}\right)^{1/(p-1)} y\left(\sqrt{\frac{q(r)}{V(r)}}t\right)$$

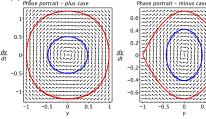
$$\ddot{y} + y \pm |y|^{p-1}y = 0$$

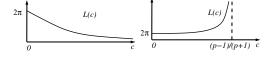
periodic orbits y(t; c), period L(c)

- c= value of first integral
- How to choose c = c(r)?



by assumptions ⇒ result





(B) Scalar breather examples via calc.var.

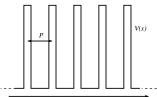


Ansatz:

$$E(x,t) = \begin{pmatrix} 0 \\ 0 \\ u(x_1,t) \end{pmatrix} \quad \hookrightarrow \quad \nabla \times \nabla \times E(x,t) = \begin{pmatrix} 0 \\ 0 \\ -\partial_{x_1}^2 u(x_1,t) \end{pmatrix}$$

Scalar semilinear wave equation:

$$(*) \begin{cases} V(x)u_{tt} - u_{xx} = \Gamma(x)|u|^{p-1}u \text{ in } \mathbb{R} \times \mathbb{R} \\ u(x,t) \to 0 \text{ as } |x| \to \infty \\ u(x,t+T) = u(x,t) \end{cases}$$



Two approaches via calc.var.:

- A. Hirsch & W.R.: $\Gamma \equiv \gamma$, $1 < \rho < \frac{5}{3}$, $V(x) = \alpha + \beta \delta^{per,P}(x)$, $\beta = \frac{8\alpha P}{\pi}$
- $\Gamma(x) = \gamma \delta_0(x), p = 3, V(x) = \alpha + \beta \delta^{per,P}(x), \beta > \frac{4\alpha P}{\pi}$ W.R.:

(B) Scalar breather examples via calc.var.



(*)
$$\begin{cases} V(x)u_{tt} - u_{xx} &= \Gamma(x)|u|^{p-1}u \text{ in } \mathbb{R} \times \mathbb{R} \\ u(x,t) &\to 0 \text{ as } |x| \to \infty \\ u(x,t+T) &= u(x,t) \end{cases}$$

with
$$V(x) = \alpha + \beta \delta^{per,P}$$
, $\alpha, \beta > 0$

(B) Scalar breather examples via calc.var.



(*)
$$\begin{cases} V(x)u_{tt} - u_{xx} &= \Gamma(x)|u|^{p-1}u \text{ in } \mathbb{R} \times \mathbb{R} \\ u(x,t) &\to 0 \text{ as } |x| \to \infty \\ u(x,t+T) &= u(x,t) \end{cases}$$

with
$$V(x) = \alpha + \beta \delta^{per,P}$$
, $\alpha, \beta > 0$

Common features:

- lacksquare $\delta^{per,P}$ is the P-periodic extension of the $\delta_{P/2}$ -distribution on x-axis
- time-period $T = 4P \sqrt{\alpha}$
- Fourier-decomp. of solution $u(x,t) = \sum_{k \text{ odd}} u_k(x)e^{ik\omega t}$, $u_{-k} = \bar{u}_k$.
- Fourier-decomp. of wave operator $L = V(x)\partial_t^2 \partial_x^2$:

$$\sigma(L) = \bigcup_{k \text{ odd}} \sigma(L_k) = \bigcup_{k \text{ odd}} \sigma(-\partial_x^2 - k^2 \omega^2 V(x))$$

• choice of α , P, β > $4\alpha P/\pi \Rightarrow 0 \notin \sigma(L)$



$$\begin{cases} V(x)u_{tt} - u_{xx} &= \gamma \delta_0(x)u^3 \text{ in } \mathbb{R} \times \mathbb{R} \\ u(x,t) &\to 0 \text{ as } |x| \to \infty \\ u(x,t+T) &= u(x,t) \end{cases}$$



Even solutions $u(x, t) = u(-x, t) \Rightarrow$ nonlinear Neumann problem

(*)
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Even solutions $u(x, t) = u(-x, t) \Rightarrow$ nonlinear Neumann problem

$$(nN) \qquad \begin{cases} V(x)u_{tt} - u_{xx} &= 0 \text{ in } (0, \infty) \times \mathbb{R}, \\ -2u_x(0, t) &= \gamma u(0, t)^3, \\ u(x, t) &\to 0 \text{ as } x \to \infty \\ u(x, t + T) &= u(x, t) \end{cases}$$



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$$\begin{cases} V(x)u_{tt} - u_{xx} &= 0 \text{ in } (0, \infty) \times \mathbb{R}, \\ -2u_{x}(0, t) &= \gamma u(0, t)^{3}, \\ u(x, t) &\to 0 \text{ as } x \to \infty \\ u(x, t + T) &= u(x, t) \end{cases}$$

Theorem 2 (R. 2016)

Let $V(x) = \alpha + \beta \delta^{per,P}(x)$ where $\alpha, P > 0$, $\beta > 4\alpha P/\pi$ and $\gamma \neq 0$. Then \exists a real-valued breather which is even in x, T/2-antiperiodic in t with $T = 4P\sqrt{\alpha}$.



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Steps:

- $u(x,t) = \sum_{k \text{ odd}} a_k \phi_k(x) e^{ik\omega t}$
- $\phi_k = \text{Bloch-mode}$ $L_k \phi_k = 0 \text{ on } (0, \infty), \text{ exp. decaying at } +\infty, \phi_k(0) = 1$
- variational problem for coefficients $(a_k)_{kodd}$, $a_{-k} = \bar{a}_k$



$$(nN) \qquad \begin{cases} V(x)u_{tt} - u_{xx} &= 0 \text{ in } (0, \infty) \times \mathbb{R}, \\ -2u_x(0, t) &= \gamma u(0, t)^3, \\ u(x, t) &\to 0 \text{ as } x \to \infty \\ u(x, t + T) &= u(x, t) \end{cases}$$

with Fourier-Bloch-decomposition $u(x,t) = \sum_{k \text{odd}} a_k \phi_k(x) e^{ik\omega t}$



$$(nN) \begin{cases} V(x)u_{tt} - u_{xx} &= 0 \text{ in } (0, \infty) \times \mathbb{R}, \\ -2u_x(0, t) &= \gamma u(0, t)^3, \\ u(x, t) &\to 0 \text{ as } x \to \infty \\ u(x, t + T) &= u(x, t) \end{cases}$$

with Fourier-Bloch-decomposition
$$u(x,t) = \sum_{k \text{odd}} a_k \phi_k(x) e^{ik\omega t}$$

$$u_X(0,t) = \sum_{k \text{odd}} \phi_k'(0) a_k e^{ik\omega t}, \quad u(0,t)^3 = \sum_{k \text{odd}} (a*a*\bar{a})_k e^{ik\omega t}$$



$$(nN) \qquad \begin{cases} V(x)u_{tt} - u_{xx} &= 0 \text{ in } (0, \infty) \times \mathbb{R}, \\ -2u_x(0, t) &= \gamma u(0, t)^3, \\ u(x, t) &\to 0 \text{ as } x \to \infty \\ u(x, t + T) &= u(x, t) \end{cases}$$

with Fourier-Bloch-decomposition $u(x,t) = \sum_{k \text{ odd}} a_k \phi_k(x) e^{ik\omega t}$

$$u_{x}(0,t) = \sum_{k \text{ odd}} \phi_{k}'(0) a_{k} e^{ik\omega t}, \quad u(0,t)^{3} = \sum_{k \text{ odd}} (a*a*\bar{a})_{k} e^{ik\omega t}$$

u weakly solves (nN) \Leftrightarrow J'[a] = 0, $a \in H$ with

$$J[a] := \sum_{k=2l+1} \underbrace{\phi_k'(0)}_{\approx -(-1)^l |k|} |a_k|^2 - \frac{\gamma}{4} |(a*a)_k|^2, \ \ H := \left\{ a_{-k} = \bar{a}_k, \sum_{k \text{ odd}} |k| |a_k|^2 < \infty \right\}$$



$$\begin{cases} V(x)u_{tt} - u_{xx} &= 0 \text{ in } (0, \infty) \times \mathbb{R}, \\ -2u_x(0, t) &= \gamma u(0, t)^3, \\ u(x, t) &\to 0 \text{ as } x \to \infty \\ u(x, t + T) &= u(x, t) \end{cases}$$

with Fourier-Bloch-decomposition $u(x,t) = \sum_{k \text{ odd}} a_k \phi_k(x) e^{ik\omega t}$

$$u_{\scriptscriptstyle X}(0,t) = \sum_{k {
m odd}} \phi_k'(0) a_k e^{ik\omega t}, \quad u(0,t)^3 = \sum_{k {
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abstract critical point theorem (Szulkin-Weth 2010):

 $\Rightarrow \exists$ truly polychromatic ground state

(B) Summary on semilinear wave-equations



Polychromatic waves $E(x,t) = \sum_k U_k(x)e^{ik\omega t}$, $U_k = \bar{U}_{-k}$ for

(semi)
$$\nabla \times \nabla \times E + V(x)\partial_t^2 E = \Gamma(x)|E|^{p-1}E$$

- vector case: radial symmetry, $E(x,t) = \psi(|x|,t) \frac{x}{|x|}$, ODE in time
- scalar case: p = 3, V(x) = cst.+periodic delta, $\Gamma(x) = \gamma$ delta at 0
 - \bullet 0 $\notin \sigma(-\partial_x^2 + V(x)\partial_t^2)$
 - use Fourier-Bloch decomposition $\sum_{k \text{ odd}} a_k \phi_k(x) e^{ik\omega t}$
 - indefinite variational problem for $(a_k)_{kodd}$
 - solve by abstract critical point theorem of Szulkin, Weth 2010
- **scalar case:** 1 , <math>V(x) = cst.+periodic delta, $\Gamma(x) = \text{cst.}$
 - -> details in talk by Andreas Hirsch, Tuesday, 11:25



$$\nabla \times \nabla \times E + \partial_t^2 (V(x)E + f(x, |E|^2)E) = 0$$

Our approaches

monochrom. \mathbb{C} -valued waves: $E(x,t) = U(x)e^{i\omega t}$

$$\nabla \times \nabla \times U + \tilde{V}(x)U = \tilde{f}(x, |U|^2)U$$
 in \mathbb{R}^3

with
$$\tilde{V} = -\omega^2 V$$
, $\tilde{f} = \omega^2 f$. elliptic, variational



$$\nabla \times \nabla \times E + \partial_t^2 \Big(V(x)E + f(x, |E|^2)E \Big) = 0$$

Our approaches

 \blacksquare $E(x,t) = U(x)\cos(\omega t)$ works for time-averaged material response

$$f\left(x, \frac{1}{T} \int_0^T |E|^2 dt\right) E$$



$$\nabla \times \nabla \times E + \partial_t^2 (V(x)E + f(x, |E|^2)E) = 0$$

Our approaches

monochrom. \mathbb{C} -valued waves: $E(x,t) = U(x)e^{i\omega t}$

$$\nabla \times \nabla \times U + \tilde{V}(x)U = \tilde{f}(x, |U|^2)U$$
 in \mathbb{R}^3

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$$\tilde{V} = -\omega^2 V$$
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$$\nabla \times \nabla \times E + \partial_t^2 \Big(V(x)E + f(x, |E|^2)E \Big) = 0$$

Our approaches

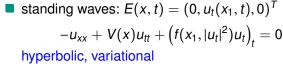
- monochrom. \mathbb{C} -valued waves: $E(x,t) = U(x)e^{i\omega t}$ $\nabla \times \nabla \times U + \tilde{V}(x)U = \tilde{f}(x, |U|^2)U$ in \mathbb{R}^3 with $\tilde{V} = -\omega^2 V$, $\tilde{f} = \omega^2 f$, elliptic, variational
- propagation direction traveling waves: $E(x, t) = (0, u(x_1 - \sqrt{\lambda}t, x_3), 0)^T$ $(-\Delta_{x_1,x_3} + \lambda V(x_3))u + \lambda (f(x_3,|u|^2)u)_{x_1x_1} = 0$ hyperbolic, bifuraction w.r.t. λ from guided modes. Talk by Piotr Idzik, Tuesday, 11:50

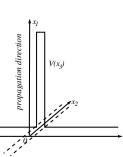


$$\nabla \times \nabla \times E + \partial_t^2 \Big(V(x)E + f(x, |E|^2)E \Big) = 0$$

Our approaches

- monochrom. \mathbb{C} -valued waves: $E(x,t) = U(x)e^{i\omega t}$ $\nabla \times \nabla \times U + \tilde{V}(x)U = \tilde{f}(x,|U|^2)U \text{ in } \mathbb{R}^3$ with $\tilde{V} = -\omega^2 V$, $\tilde{f} = \omega^2 f$. elliptic, variational
- traveling waves: $E(x, t) = (0, u(x_1 \sqrt{\lambda}t, x_3), 0)^T$ $(-\Delta_{x_1, x_3} + \lambda V(x_3))u + \lambda (f(x_3, |u|^2)u)_{x_1x_1} = 0$ hyperbolic, bifuraction w.r.t. λ from guided modes. Talk by Piotr Idzik, Tuesday, 11:50





(C) Elliptic Curl-Curl problem



$$\nabla \times \nabla \times U + V(x)U = f(x, |U|^2)U$$
 in \mathbb{R}^3

- (0) $U(x_1, x_2, x_3) = (0, 0, u(x_1, x_2))^T$ leads to NLS (many results!)
- (1) Benci-Fort. ('04) & Azzollini-B.-d'Aprile-F. ('06) & d'A.-Siciliano ('11), Zeng ('16):

$$\nabla \times \nabla \times U = f(|U|^2)U$$
 in \mathbb{R}^3

Existence of ground-states in subspaces of cylindrical symmetry

(2) Bartsch-Mederski ('15,'17), survey ('17 J. Fixed Point Theory Appl.):

$$\nabla \times \mu(x)^{-1} \nabla \times U - \omega^2 \epsilon(x) U = \partial_U F(x, U) \text{ in } \Omega, \quad v \times U = 0 \text{ on } \partial\Omega.$$

(3) Mederski('15): $f(s) \approx |s|^{\frac{p-1}{2}}$ near $0, f(s) \approx |s|^{\frac{q-1}{2}}$ near $\infty, 1 .$

$$\nabla \times \nabla \times U + V(x)U = f(|U|^2)U$$
 in \mathbb{R}^3

- Mederski ('16): Brezis-Nirenberg Curl-Curl problem: Tuesday, 14:00
- Bartsch-Dohnal-Plum-R. ('14) & Hirsch-R. ('16) ... next

(C) Common variational set-up



$$J[U] = \int_{\mathbb{R}^3} |\nabla \times U|^2 + \tilde{V}(x)|U|^2 - \tilde{F}(r,z,|U|^2) dx,$$

$$U \in X = H(\operatorname{curl}; \mathbb{R}^3) \cap L^{p+1}(\mathbb{R}^3)$$

Here is the problem: $\|\nabla U\|_{L^{2}}^{2} = \|\nabla \times U\|_{L^{2}}^{2} + \|\nabla \cdot U\|_{L^{2}}^{2}$.

Constraint $\{U : \text{div } U = 0\}$ does not solve it \Rightarrow Lagrange-multiplier!

$$U(r,z) := u(r,z) \begin{pmatrix} -x_2 \\ x_1 \\ 0 \end{pmatrix} \implies \text{div } U = 0.$$

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Symmetries! Look for cylindrical symmetry in coordinates (r, z):

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$$-\Delta_5 u(r,z) + \tilde{V}(r,z)u = \tilde{f}(r,z,r^2u^2)u$$
 for $r > 0, z \in \mathbb{R}$.

This is a NLS-type equation in \mathbb{R}^5 !

(C) Results I - (Bartsch-Dohnal-Plum-R., NoDeA 2016)



(*)
$$\nabla \times \nabla \times U + \tilde{V}(x)U = \tilde{\Gamma}(x)|U|^{p-1}U \quad \text{in} \quad \mathbb{R}^3$$

General assumption:
$$\tilde{V} = \tilde{V}(r, x_3), \tilde{\Gamma} = \tilde{\Gamma}(r, x_3), r = \sqrt{x_1^2 + x_2^2}$$

Theorem 3 (Defocusing case)

- $\tilde{\Gamma}(x) \leq -C(1+|x|^{\alpha}), \, \alpha > \frac{3}{2}(p-1), \, p > 1,$
- $\quad \blacksquare \ \, \tilde{V} \in L^{\infty}(\mathbb{R}^3), \, \sup \, \tilde{V} < 0.$

Then (*) has a (restricted) ground-state.

Theorem 4 (Focusing case)

- inf $\tilde{\Gamma} > 0$, \tilde{V} , $\tilde{\Gamma} \in L^{\infty}(\mathbb{R}^3)$ are 1-periodic in x_3 ,
- $\blacksquare 1$
- \blacksquare 0 $\notin \sigma(L)$ with $L = \nabla \times \nabla \times + \tilde{V}(x)$.

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Theorem 5 (Positive definite case, focusing)

- \bullet 0 < min $\sigma(\nabla \times \nabla \times + \tilde{V})$
- $0 \le \tilde{f}(r,z,s) \le C(1+s^{\frac{p-1}{2}}), 1$
- $\tilde{f}(r,z,s) = o(1)$ as $s \to 0$ uniformly in r, z
- \blacksquare $s \mapsto \tilde{f}(r,z,s)$ strictly increasing in s
- $\tilde{F}(r,z,s)/s \to \infty$ as $s \to \infty$ uniformly in r, z
- $\tilde{V}(r,z)$ reverse Steiner-symmetric in z
- $\phi_{\sigma}(r,z,s) := \tilde{f}(r,z,(s+\sigma)^2)(s+\sigma)^2 \tilde{f}(r,z,s^2)s^2$ is symmetrically decreasing in z for all $s \ge 0$, $\sigma \ge 0$

Ex.:
$$\tilde{f}(r, z, |U|^2) = \tilde{\Gamma}(r, z) \log(1 + |U|^2), = \tilde{\Gamma}(r, z) |U|^{p(z)-1}, \overline{\text{Rg}(p)} \subset (1, 5)$$



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(C) Sketch of variational existence proof



$$-\Delta_5 u(r,z) + V(r,z)u = f(r,z,r^2u^2)u \text{ for } r > 0, z \in \mathbb{R}.$$

$$J[u] = \int_{\mathbb{R}^5} |\nabla u|^2 + V(r,z)u^2 - \frac{F(r,z,r^2u^2)}{r^2} dx^5, \quad u \in H^1_{cyl}(\mathbb{R}^5)$$

Minimize J over the Nehari-manifold [cf. Szulkin-Weth, '10]:

$$\mathcal{N} = \left\{ u \neq 0; \, N[u] = \int_{\mathbb{R}^5} |\nabla u|^2 + V(r, z)u^2 - f(r, z, r^2u^2)u^2 \, dx^5 = 0 \right\}$$

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Changing from u to $|u|^*$ (Steiner symmetrization w.r.t. z) we get

$$J[|u|^*] \le J[u], \quad N[|u|^*] \le N[u]$$

because of the condition [Brock, '00]

$$\phi_{\sigma}(r,z,s) := f(r,z,(s+\sigma)^2)(s+\sigma)^2 - f(r,z,s^2)s^2$$
 symm z.

Moreover: weak sequ. cont.'y along $(|u_k|^*)_{k\in\mathbb{N}}$ [inspired by Lions,'81,'82].



$$\begin{cases} V(x)u_{tt} - u_{xx} &= \gamma \delta_0(x)(u_t^3)_t \text{ in } \mathbb{R} \times \mathbb{R} \\ u(x,t) &\to 0 \text{ as } |x| \to \infty \\ u(x,t+T) &= u(x,t) \end{cases}$$



Even solutions $u(x,t) = u(-x,t) \Rightarrow$ nonlinear Neumann problem

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Theorem 6 (R. 2016)

Let $V(x) = \alpha + \beta \delta^{per,P}(x)$ where $\alpha, P > 0, \beta > 4\alpha P/\pi$ and $\gamma \neq 0$. Then \exists a real-valued breather which is even in x, T/2-antiperiodic in t with $T=4P\sqrt{\alpha}$.



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Steps:

- $u(x,t) = \sum_{k \text{ odd}} \frac{a_k}{L} \phi_k(x) e^{ik\omega t}, \phi_k = \text{normalized Bloch-mode}$
- variational problem for coefficients $(a_k)_{kodd}$, $a_{-k} = -\bar{a}_k$



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with Fourier-Bloch-decomposition $u(x,t) = \sum_{k \text{odd}} \frac{a_k}{k} \phi_k(x) e^{ik\omega t}$



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u weakly solves (nN) \Leftrightarrow J'[a] = 0, $a \in H$ with

$$J[a] := \sum_{k=2l+1} \frac{\omega^4}{4} |(a*a)_k|^2 + \underbrace{\phi'_k(0)}_{\approx -(-1)^l |k|} \frac{|a_k|^2}{k^2 \gamma}, \ \ H := \left\{ a_{-k} = -\bar{a}_k, \|a*a\|_{l^2} < \infty \right\}$$



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minimizer = polychromatic ground state

(C) Summary on quasilinear wave-equations



(quasi)
$$\nabla \times \nabla \times E + \partial_t^2 (V(x)E + \Gamma(x)|E|^{p-1}E) = 0$$

- ∃ monochromatic ground-state $E(x,t) = U(x)e^{i\omega t}$:
 - elliptic vector case in \mathbb{R}^3 with cylindrical symmetry
 - defocusing case: $\Gamma(x) \leq -C(1+|x|^{\alpha}), \alpha > \frac{3}{2}(p-1), \text{ inf } V > 0$
 - focusing case: periodic structure in z, $0 \notin \sigma(\nabla \times \nabla \times -\omega^2 V(x))$
 - focusing case: Steiner symmetry in z, $0 < \sigma(\nabla \times \nabla \times -\omega^2 V(x))$

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- \blacksquare \exists scalar 1+1-dim \mathbb{R} -valued standing polychromatic waves:
 - p = 3, V(x) = cst.+periodic delta, $\Gamma(x) = \gamma$ delta at 0
 - $0 \notin \sigma(-\partial_x^2 + V(x)\partial_t^2)$
 - use Fourier-Bloch decomposition $\sum_{k \text{odd}} \frac{a_k}{k} \phi_k(x) e^{ik\omega t}$
 - coercive variational problem for $(a_k)_{kodd} \hookrightarrow minimizer=breather$

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- \blacksquare \exists scalar 2+1-dim \mathbb{R} -valued traveling waves:
 - $p=3,\ V(x_3)=$ delta at 0, $\Gamma(x_3)\in L^\infty$: use bifurcation theory
 - -> details in talk by Piotr Idzik, Tuesday, 11:50